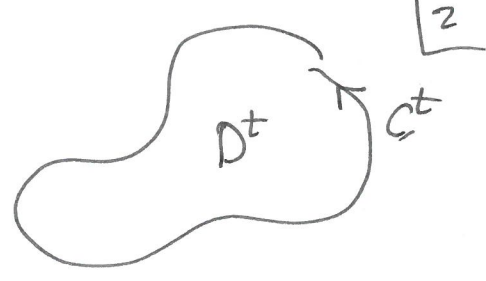


DAY 04



STILL FOCUSED ON THE INCOMP. EULER EQNS W/ CONST DENSITY:

$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u = -\nabla p \\ \nabla \cdot u = 0 \end{cases}$$

~~external forces~~

★ TODAY EVERYTHING WILL BE IN $n=2$ DIMENSIONS.

Last time we saw that the assumptions $w = \nabla \times u \equiv 0$, and $\frac{\partial u}{\partial t} \equiv 0$ simplify things considerably. Today we will drop these assumptions and look at the Cauchy problem

$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u = -\nabla p \\ \nabla \cdot u = 0 \\ u|_{t=0} = u_0 \end{cases}$$

The vorticity w will play a prominent role. Cauchy data

WHAT ARE SOME MORE FAMILIAR CAUCHY PROBLEMS?

The vorticity eqn in 2D

Consider a domain $D^t = \mathbb{R}^2(0^t)$ with boundary $C^t = \partial D^t = \mathbb{R}^1(C^0)$.
By Kelvin's Circulation theorem,

$$\frac{d}{dt} \int_{C^t} u \cdot dl = 0.$$

On the other hand by Green's theorem

$$\int_{C^t} u \cdot dl = \int_{D^t} w \, dx$$

Thus $\frac{d}{dt} \int_{D^t} w \, dx = 0.$

on the other hand we have our usual chain rule formula:

$$\frac{d}{dt} \int_{D^t} w \, dx = \int_{D^t} \left(\frac{Dw}{Dt} + w(\nabla \cdot u) \right) dx$$

and so we ~~we~~ get $\int_{D^t} \frac{Dw}{Dt} \, dx = 0.$

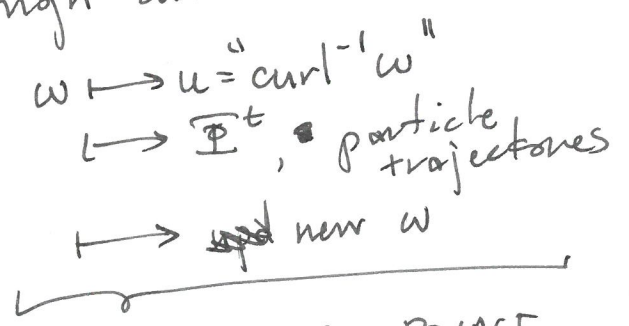
Since D^t was arbitrary,

$$\frac{Dw}{Dt} = \frac{\partial w}{\partial t} + (u \cdot \nabla)w = 0$$

This means that

w is constant along particle trajectories

This is ~~extremely~~ extremely useful and only valid in 2D.
Many arguments follow the rough outline:



FIND A FIXED POINT

3] Some background

Fixed-point theorems

X is a Banach space

Banach: If $F: X \rightarrow X$ has $|F(x) - F(y)| < \rho|x - y|$ for some fixed $\rho < 1$, then F has a unique fixed point.
 unit ball $\subseteq \mathbb{R}^n$

Brouwer If $f: B \rightarrow B$ is continuous then it has at least one fixed point.

Schauder If $K \subseteq X$ is convex and compact, then every ~~cont~~ continuous $F: K \rightarrow K$ has at least one fixed point.

Idea: if K were finite dimensional then this is Brouwer. But compactness means it is "almost" finite dimensional.

Recovering u from w

Since we're in 2D and $\nabla \cdot u = 0$, we know that

$$u = \nabla^\perp \psi =: (\partial_y \psi, -\partial_x \psi)$$

for some stream function ψ .

So

$$\begin{aligned} w = \nabla \times u &= -\nabla^\perp \cdot u \\ &= -\nabla^\perp \cdot \nabla^\perp \psi \\ &= -\Delta \psi. \end{aligned}$$

If we could invert $-\Delta$, then this would give

$$\psi = (-\Delta)^{-1} w \quad [4]$$

$$u = \nabla^\perp \psi = \nabla^\perp (-\Delta)^{-1} w$$

This brings us to

Basic potential theory (2D!)

If a ^{smooth} function f has compact support \checkmark then in \mathbb{R}^2

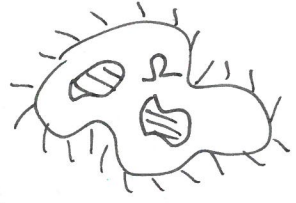
$$f(x) = \int_{\mathbb{R}^2} \frac{1}{2\pi} \log|x-y| \Delta f(y) dy$$

$$= \frac{1}{2\pi} \log|\cdot| * \Delta f$$

"fundamental solution"

For a smooth bold domain Ω , the problem

$$\begin{cases} \Delta f = g & \text{in } \Omega \\ f = 0 & \text{on } \partial\Omega \end{cases}$$



similarly has a unique soln, expressible as

$$f(x) = \int_{\Omega} G^{\Omega}(x,y) g(y) dy$$

(*)

Green's function for $-\Delta$ on Ω

G^{Ω} is smooth away from $x=y$, and near $x=y$ it looks like

$$G^{\Omega}(x,y) = -\frac{1}{2\pi} \log|x-y| + \text{smooth terms}$$

We write (*) as

$$f = (-\Delta)^{-1} g.$$

5] Hölder spaces

Unfortunately, while $u \in C^{k+2} \Rightarrow \Delta u \in C^k$ for $k=0, 1, 2, \dots$, the reverse implication is NOT TRUE. In particular, $u = (x_1^2 - x_2^2)(-\log|x|)^{1/2} \notin C^2$

while $\Delta u = \text{ugly} \dots \in C^0$. So we need different spaces of functions which somehow have "fractional" regularity. *big list/tensor of all k-th order partials*

Def If $u, Du, \dots, D^k u$ exist and are continuous on Ω we say $u \in C^k(\Omega)$. If they continuously extend to $\bar{\Omega}$ then we say $u \in C^k(\bar{\Omega})$. $C^k(\bar{\Omega})$ is a Banach space with the norm

$$|u|_{C^k(\bar{\Omega})} = \sum_{j=0}^k \sup_{\bar{\Omega}} |D^j u|$$

Hölder semi-norm

Def For $0 < \theta < 1$ we set $[u]_{C^\theta(\bar{\Omega})} = \sup_{\substack{x, y \in \bar{\Omega} \\ x \neq y}} \frac{|u(x) - u(y)|}{|x - y|^\theta}$

When $[u]_{C^\theta} < \infty$ we say that u is Hölder continuous with exponent θ . **Ex: $u(x) = |x|^\theta$**

Def For $k=0, 1, 2, \dots$ and $\theta \in (0, 1)$ we set

$$|u|_{C^{k+\theta}} = |u|_{C^k} + [D^k u]_{C^\theta}$$

The corresponding Banach space is called $C^{k+\theta}(\bar{\Omega})$. **EX: $|x|^{k+\theta}, x \geq 0$**

Fact For any $f \in C^{k+\theta}(\bar{\Omega})$, the problem

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

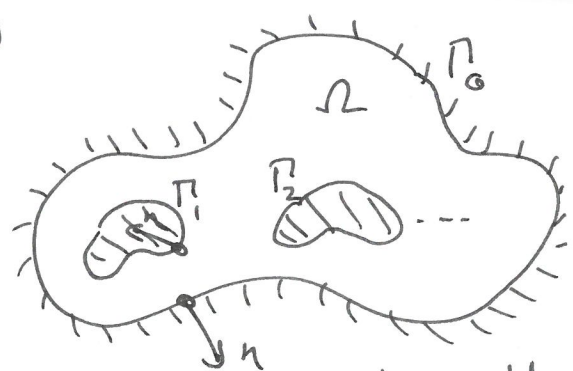
has a unique solution $u \in C^{k+2+\theta}(\bar{\Omega})$. The corresponding linear map $(-\Delta)^{-1}: u \mapsto f$ is bold $C^{k+\theta}(\bar{\Omega}) \rightarrow C^{k+2+\theta}(\bar{\Omega})$

(Gilbarg-Trudinger Cor 4.14)

SO, these are good spaces to study $(-\Delta)^{-1}$ in, and hence good spaces to use when we want to use formulas like $u = \nabla^\perp (-\Delta)^{-1} w$

• Kato 1967: On classical solutions of the Two-Dim. Non-stationary Euler Eqn

This paper is available online and is very readable. [Some of the references are bit old, and some of the notation is less fashionable] We will change notation in a few places to match what we've been using so far.



Let Ω be a smooth bdd domain $\subseteq \mathbb{R}^2$, with boundary components $\Gamma_0, \Gamma_1, \dots, \Gamma_m$ as above. Consider the Euler equation in $\bar{Q}_T = \Omega \times [0, T]$

(E)
$$\begin{cases} \frac{\partial u}{\partial t} + (u \cdot \nabla)u = -\nabla p + f \\ \nabla \cdot u = 0 \\ u \cdot n = 0 \text{ on } \partial\Omega \\ u|_{t=0} = u_0 \end{cases}$$

Then (Kato) Suppose $u_0 \in C^{1+\theta}(\bar{\Omega})$, $\nabla \cdot u_0 = 0$, $a \cdot n = 0$ on Γ , and $f \in C^{1+\theta, 0}(\bar{Q}_T)$ for some Hölder exponent $\theta \in (0, 1)$. Then \exists a soln (u, p) of (E) such that u, p & all their derivatives appearing in (E) are $C(\bar{Q}_T)$. Such a solution is unique up to an arbitrary function of t which may be added to p .

Simplifications we will assume that $\Gamma = \Gamma_0$ so Ω is simply-connected $f = 0$ so no external forces

Also we will not completely prove every lemma, "Exercise for the reader"

Definitions and notation

$\nabla \cdot u = 0 \iff$ "u is a flow"
 $\left. \begin{matrix} \nabla \cdot u = 0, \\ u \cdot n = 0 \text{ on } \Gamma \end{matrix} \right\} \iff$ "u is a tangential flow"

$(u, v) = \int_{\Omega} u \cdot v \, dx = (u, v)_{L^2}$

$|u|_{L^2}^2 = (u, u)$

$u \in C^k$ always means $u \in C^k(\bar{\Omega})$ or $C^k(\bar{Q}_T)$ and never just $u \in C^k(\Omega)$.

$u \in C^{j,k}(\bar{Q}_T)$ means

$D_x^p D_t^q u \in C^0(\bar{Q}_T)$

for $0 \leq p \leq j, 0 \leq q \leq k$.

For $0 \leq \delta < 1, 0 \leq \varepsilon < 1$, $u \in C^{j+\delta, k+\varepsilon}(\bar{Q}_T)$ means that $u \in C^{j,k}(\bar{Q}_T)$ and that

$[D_x^p D_t^q u]_{C^{\delta, \varepsilon}} < \infty$

for $0 \leq p \leq j, 0 \leq q \leq k$, where

$[u]_{C^{\delta, \varepsilon}} := \frac{1}{\delta} \sup_{s>0} \sup_{x \neq y} \frac{|u(x, t) - u(y, t)|}{|x - y|^\delta} + \frac{1}{\varepsilon} \sup_{x} \sup_{t \neq s} \frac{|u(x, s) - u(x, t)|}{|x - t|^\varepsilon}$

10] What we know is that

§ I Preliminaries

$$[\varphi]_{C^{\delta, \varepsilon}(\bar{Q}_T)} = \sup_t \sup_{x \neq y} \frac{|\varphi(x, t) - \varphi(y, t)|}{|x - y|^\delta} + \sup_x \sup_{t \neq s} \frac{|\varphi(x, s) - \varphi(x, t)|}{|s - t|^\varepsilon}$$

~~...~~
∞

$$|\varphi|_{C^0(\bar{Q}_T)} = \sup_{x, t} |\varphi| < \infty$$

The first term in $\textcircled{*}$ is ~~easy~~ easy to bound:

$$\begin{aligned} & \sup_t |\varphi(\cdot, t)|_{C^{\lambda \delta}} \\ &= \sup_t \left(\sup_x |\varphi| + \sup_{x \neq y} \frac{|\varphi(x, t) - \varphi(y, t)|}{|x - y|^{\lambda \delta}} \right) \\ &\leq |\varphi|_{C^0(\bar{Q}_T)} + (\text{diam } \Omega) [\varphi]_{C^{\delta, \varepsilon}(\bar{Q}_T)} \end{aligned}$$

The next term is a bit more complicated:

$$\begin{aligned} & \sup_{t \neq s} \frac{|\varphi(\cdot, t) - \varphi(\cdot, s)|}{|t - s|^{(1-\lambda)\varepsilon}} \\ &= \sup_{t \neq s} \sup_x \frac{|\varphi(x, t) - \varphi(x, s)|}{|t - s|^{(1-\lambda)\varepsilon}} \\ & \quad + \sup_{t \neq s} \sup_{x \neq y} \frac{|\varphi(x, t) - \varphi(y, t) - (\varphi(y, t) - \varphi(y, s))|}{|x - y|^{\lambda \delta} |t - s|^{(1-\lambda)\varepsilon}} \\ &\leq T^{\lambda \varepsilon} [\varphi]_{C^{\delta, \varepsilon}} + [\varphi]_{C^{\delta, \varepsilon}}^{\lambda} [\varphi]_{C^{\delta, \varepsilon}}^{1-\lambda} \end{aligned}$$

(ii) Exercise

Lemma 1.1 Let $u \in C^1(\bar{\Omega})$ be a flow and set $w = \nabla_x u = -\nabla^\perp u$. Then

$$(\varphi, u \cdot \nabla \Phi) + (u \cdot \nabla) u, \nabla^\perp \Phi = 0$$

for all $\Phi \in C^1(\bar{\Omega})$ which vanish on Γ .

Sketch
Pf: When $u \in C^2(\bar{\Omega})$, this can be proven by integrating by parts (divergence theorem and/or Green's theorem).

For $u \in C^1(\bar{\Omega})$, first extend to $\tilde{u} \in C^1(\tilde{\Omega}^*)$ where $\tilde{\Omega}^* \supset \bar{\Omega}$ and then apply to

$$u^\varepsilon = \phi_\varepsilon * u \in C^\infty$$

sequence of mollifiers approximating δ

Lemma 1.2 (i) If $\varphi \in C^{\delta, \varepsilon}(\bar{Q}_T)$ with $\delta, \varepsilon \in (0, 1)$, then

$$t \mapsto \varphi(\cdot, t)$$

is a $C^{(1-\lambda)\varepsilon}$ function $[0, T] \rightarrow C^{\lambda \delta}(\bar{\Omega})$

for any $\lambda \in (0, 1)$. (ii) If $\varphi \in C^{\delta, 0}(\bar{Q}_T)$ with $\delta \in (0, 1)$ then it is a uniformly continuous map $[0, T] \rightarrow C^{\delta'}(\bar{\Omega})$

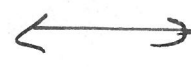
for any $\delta' < \delta$.

Pf (ii) ~~...~~ We need to show that

$$\begin{aligned} & |\varphi|_{C^{(1-\lambda)\varepsilon}([0, T]; C^{\lambda \delta}(\bar{\Omega}))} \\ &= \sup_t |\varphi(\cdot, t)|_{C^{\lambda \delta}(\bar{\Omega})} \\ & \quad + \sup_{t \neq s} \frac{|\varphi(\cdot, t) - \varphi(\cdot, s)|_{C^{\lambda \delta}}}{|t - s|^{(1-\lambda)\varepsilon}} < \infty \end{aligned}$$

$\textcircled{*}$

12] OOPS, ~~COLUMNS~~ COLUMNS REVERSED



Lem 1.3 Let $S \in (0, 1)$. Then 11

What if we just have $\varphi \in L^\infty$?
Then we have to directly use the Green's representation

$$u^{(x)} = \int \nabla_x^\perp G^{\mathbb{R}^2}(x, y) \varphi(y) dy$$

Now from

$$G^{\mathbb{R}^2}(x, y) = -\frac{1}{2\pi} \log|x-y| + \text{smooth}$$

we get

$$|D_x G^{\mathbb{R}^2}(x, y)| \leq \frac{K}{|x-y|}$$

$$|D_x^2 G^{\mathbb{R}^2}(x, y)| \leq \frac{K}{|x-y|^2}$$

for some constant K .

Lem 1.4 If $\varphi \in L^\infty(\mathbb{R}^2)$, then

$(-\Delta)^{-1} \varphi \in C^2(\mathbb{R}^2)$ and

$u = -\nabla^\perp (-\Delta)^{-1} \varphi$ is a tangential flow with

$$\|u\|_{L^\infty} \leq K \|\varphi\|_{L^\infty}$$

$$\|u(x) - u(y)\| \leq K \|\varphi\|_{L^\infty} |x-y| \chi(|x-y|)$$

where $K = K(\mathbb{R}^2)$ and

$$\chi(s) = (1 + \log(1/s)) \mathbb{1}_{0 < s < 1}$$

PROOF NEXT TIME

(i) $\nabla^\perp (-\Delta)^{-1} : C^s(\mathbb{R}^2) \rightarrow C^{1+s}(\mathbb{R}^2)$ is continuous. Moreover

(ii) $\varphi \in C^{s,0}(\overline{\mathbb{Q}_T}) \Rightarrow u = \nabla^\perp (-\Delta)^{-1} \varphi$ is $C^{1+s',0}(\overline{\mathbb{Q}_T}) \forall s' < s$.

(iii) similarly $\varphi \in C^{s,\varepsilon}(\overline{\mathbb{Q}_T}) \Rightarrow u \in C^{1+s',\varepsilon'} \forall s' < s, \varepsilon' < \varepsilon$.

Pf (i) We know $(-\Delta)^{-1} : C^s(\mathbb{R}^2) \rightarrow C^{2+s}(\mathbb{R}^2)$ is bdd, and similarly

$$\frac{\partial}{\partial x_i} : C^{2+s}(\mathbb{R}^2) \rightarrow C^{1+s}(\mathbb{R}^2)$$

(iii) By Lem 1.2,

$$\varphi \in C^{s,\varepsilon}(\overline{\mathbb{Q}_T}) \Rightarrow \varphi \in C^{(1-\lambda)\varepsilon}([0,T], C^{\lambda s}(\mathbb{R}^2))$$

for all $\lambda \in (0, 1)$. ~~First choose λ so that $(1-\lambda)\varepsilon = \varepsilon'$. Then~~

Thus (WHY?)

$$u = \nabla^\perp (-\Delta)^{-1} \varphi \in C^{(1-\lambda)\varepsilon}([0,T], C^{1+\lambda s}(\mathbb{R}^2))$$

First choose $\lambda \in (0, 1)$ so that

$$(1-\lambda)\varepsilon = \varepsilon'. \text{ Then have } u \in C^{\varepsilon'}([0,T], C^{1+\lambda s}(\mathbb{R}^2))$$

$$\Rightarrow D_x u \in C^{\varepsilon'}([0,T], C^{\lambda s}(\mathbb{R}^2))$$

$$\Rightarrow u \in C^{0,\varepsilon'}(\overline{\mathbb{Q}_T}).$$

Similarly with $\lambda \in (0, 1)$ so that $\lambda s = s'$ we get

$$\bullet D_x u \in C^{(1-\lambda)\varepsilon}([0,T], C^{s'}(\mathbb{R}^2))$$

$$\Rightarrow u \in C^{s',0}(\overline{\mathbb{Q}_T}).$$

Thus $u \in C^{s',\varepsilon'}(\overline{\mathbb{Q}_T})$ as desired.

(ii) Exercise.